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## OPTIMIZED QUASIPARTICLE DENSITY FUNCTIONAL AND GREEN'S FUNCTIONS METHOD TO COMPUTING BOND ENERGIES OF DIATOMIC MOLECULES

It is presented an advanced approach to computing the energy and spectral parameters of the diatomic molecules, which is based on the hybrid combined density functional theory (DFT) and the Green's-functions (GF) approach. The Fermi-liquid quasiparticle version of the density functional theory is modified and used. The density of states, which describe the vibrational structure in photoelectron spectra, is defined with the use of combined DFT-GF approach and is well approximated by using only the first order coupling constants in the optimized one-quasiparticle approximation. Using the combined DFT-GF approach to computing the spectroscopic factors of diatomic molecules leads to significant simplification of the calculation procedure and increasing an accuracy of theoretical prediction. As illustration, the results of computing the bond energies in a number of known diatomic molecules are presented and compared with alternative theoretical results, obtained within discrete-variational  $X_\alpha$ , muffin-tin orbitals and other methods.

### 1. Introduction

In this paper we study the problem of calculating the important spectroscopic characteristics of multielectron systems (atoms and molecules), namely, the spectroscopic factor. The spectroscopic factor is one of the most important characteristics of atomic and molecular systems and the precise information about it is very important for many applications [1-47].

In calculations based on the density functional theory (DFT) methods in the local density (LP) approximation, invariants have become widespread: discrete-variational  $X_\alpha$  (DV- $X_\alpha$ ), muffin-tin orbitals (MTO) method in a version of the linear MTO method and localized orbitals, modified DV- $X_\alpha$  method with using a scheme of the transition state (TS) (see [1]). Although, in computational terms, these methods are highly economical, the error in calculating complex molecules based on them can reach several eV

In this paper we present an advanced approach to computing the spectroscopic factors of the diatomic molecules within the hybrid combined density functional theory (DFT) in the Fermi-liquid formulation and

the Green's-functions (GF) approach to quantitative determination of the spectroscopic factors for some molecular systems. The approach is based on the Green's function method (Cederbaum-Domske version) [1,2] and Fermi-liquid DFT formalism [3-7] and using the novel effective density functional (see also [11-22]). It is important that the calculational procedure is significantly simplified with using the quasiparticle DFT formalism.

### 2. Many-body theory

As usually, introducing a field operator  $\Psi(R, \theta, x) = \sum_i \phi_i(x, R, \theta) a_i(R, \theta)$  with the Hartree-Fock (HF) one-particle functions  $\phi_i$  ( $\epsilon_i(R)$  are the one-particle HF energies and  $f$  denotes the set of orbitals occupied in the HF ground state;  $R_0$  is the equilibrium geometry on the HF level) and dimensionless normal coordinates  $Q_s$  one can write the standard Hamiltonian as follows [2,7]:

$$H = H_E + H_N + H_{EN}^{(1)} + H_{EN}^{(2)}, \quad (1)$$

$$H_E = \sum_i \epsilon_i (R_0) a_i^\dagger a_i + \frac{1}{2} \sum V_{ijkl} (R_0) a_i^\dagger a_j^\dagger a_l a_k - \sum_{i,j} \sum_{k \in f} [V_{ikjk} (R_0) - V_{ikkj} (R_0)] a_i^\dagger a_j, \quad (2)$$

$$H_N = \hbar \sum_{s=1}^M \omega_s (b_s^\dagger b_s + \frac{1}{2}), \quad \tilde{H}_0 = \sum_{s=1}^M \hbar \omega_s b_s^\dagger b_s + \sum_{s=1}^M g_s^k (b_s + b_s^\dagger) + \sum_{s,s'=1}^M \gamma_{ss'}^k (b_s + b_s^\dagger) (b_{s'} + b_{s'}^\dagger) \quad (6)$$

$$H_{EN}^{(1)} = 2^{-1/2} \sum_{s=1}^M \left( \frac{\partial \epsilon_i}{\partial Q_s} \right)_0 (b_s + b_s^\dagger) [a_i^\dagger a_i - n_i] + \frac{1}{4} \sum_i \sum_{s,s'=1}^M \left( \frac{\partial^2 \epsilon_i}{\partial Q_s \partial Q_{s'}} \right)_0 (b_s + b_s^\dagger) (b_{s'} + b_{s'}^\dagger) [a_i^\dagger a_i - n_i] \quad (3a)$$

$$g_s^i = \pm \frac{1}{\sqrt{2}} \left( \frac{\partial \epsilon_i}{\partial Q_s} \right)_0, \quad \gamma_{ss'}^i = \pm \frac{1}{4} \left( \frac{\partial^2 \epsilon_i}{\partial Q_s \partial Q_{s'}} \right)_0. \quad (7)$$

$$H_{EN}^{(2)} = 2^{-3/2} \sum_{s=1}^M \left( \frac{\partial V_{ijkl}}{\partial Q_s} \right)_0 (b_s + b_s^\dagger) [\delta v_1 a_i^\dagger a_j^\dagger a_k a_l + \delta v_2 a_l a_k a_i^\dagger a_j^\dagger + 2 \delta v_3 a_j^\dagger a_k a_l a_i^\dagger] + \frac{1}{8} \sum_{s,s'=1}^M \left( \frac{\partial^2 V_{ijkl}}{\partial Q_s \partial Q_{s'}} \right)_0 (b_s + b_s^\dagger) (b_{s'} + b_{s'}^\dagger) \cdot [\delta v_1 a_i^\dagger a_j^\dagger a_k a_l + \delta v_2 a_l a_k a_i^\dagger a_j^\dagger + 2 \delta v_3 a_j^\dagger a_k a_l a_i^\dagger] \quad (3b)$$

with  $n_i=1$  (0),  $i \in f$  ( $i \notin f$ ),  $\delta \sigma_f=1$  (0),  $(ijkl) \in \sigma_f$ , where the index set  $v_1$  means that at least  $\phi_k$  and  $\phi_l$  or  $\phi_i$  and  $\phi_j$  are unoccupied,  $v_2$  that at most one of the orbitals is unoccupied, and  $v_3$  that  $\phi_k$  and  $\phi_j$  or  $\phi_l$  and  $\phi_i$  are unoccupied. The  $\omega_s$  are the HF frequencies;  $b_s, b_s^\dagger$  are destruction and creation operators for vibrational quanta as

$$Q_s = (1/\sqrt{2})(b_s + b_s^\dagger), \quad \partial/\partial Q_s = (1/\sqrt{2})(b_s - b_s^\dagger). \quad (4)$$

The interpretation of the above Hamiltonian and an exact solution of the one-body HF problem is given in refs. [1,-7]. The usual way is to define the HF-single-particle component  $H_0$  of the Hamiltonian (4) is as in Refs. [1,4]. Correspondingly in the one-particle picture the density of occupied states is given by

$$N_k^0(\epsilon) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} dt e^{i\hbar^{-1}(\epsilon - \epsilon_k)t} \langle 0 | e^{\pm i\hbar^{-1} \tilde{H}_0 t} | 0 \rangle, \quad (5)$$

$$\tilde{H}_0 = \sum_{s=1}^M \hbar \omega_s b_s^\dagger b_s + \sum_{s=1}^M g_s^k (b_s + b_s^\dagger) + \sum_{s,s'=1}^M \gamma_{ss'}^k (b_s + b_s^\dagger) (b_{s'} + b_{s'}^\dagger)$$

$$g_s^i = \pm \frac{1}{\sqrt{2}} \left( \frac{\partial \epsilon_i}{\partial Q_s} \right)_0, \quad \gamma_{ss'}^i = \pm \frac{1}{4} \left( \frac{\partial^2 \epsilon_i}{\partial Q_s \partial Q_{s'}} \right)_0. \quad (7)$$

To get function  $N_k(\epsilon)$  one calculates the GF  $G_{kk'}(\epsilon)$  (see details in Refs. [1-7,31-35]:

$$G_{kk'}(\epsilon) = -i\hbar^{-1} \int_{-\infty}^{\infty} dt e^{i\hbar^{-1}\epsilon t} \langle \psi_0 | T \{ a_k(t) a_k^\dagger(0) \} | \psi_0 \rangle \quad (8)$$

Choosing the unperturbed  $H_0$  to be  $H_0 = \sum \epsilon_i a_i^\dagger a_i + H_N$  one could define GF as

$$G_{kk'}^{OB}(t) = \pm \delta_{kk'} i \exp[-in^{-1}(\epsilon_k \mp \Delta \epsilon)t] \cdot \sum_n \left| \langle \hat{n}_k | U_k | 0 \rangle \right|^2 \exp(\pm in_k \cdot \hat{\omega}_k t) \quad (9)$$

The direct method for calculation of  $N_k(\epsilon)$  as the imaginary part of the GF includes a definition of the vertical I.P. (V.I.P.s) of the reference molecule and then of  $N_k(\epsilon)$ .

The zeros of the functions:

$$D_k(\epsilon) = \epsilon - [\epsilon^{op} + \Sigma(\epsilon)]_k, \quad (10)$$

where  $(\epsilon^{op} + \Sigma)_k$  denotes the  $k$ -th eigenvalue of the diagonal matrix of the one-particle energies added to matrix of the self-energy part, are the negative V. I. P. 's for a given geometry. One can write [2,4]:

$$(V.I.P.)_k = -(\epsilon_k + F_k), \quad F_k = \Sigma_{kk}(- (V.I.P.)_k) \approx \frac{1}{1 - \partial \Sigma_{kk}(\epsilon_k) / \partial \epsilon} \Sigma_{kk}(\epsilon_k) \quad (11)$$

Expanding the ionic energy  $E_k^{N-1}$  about the equilibrium geometry of the reference molecule in a power series of the normal coordinates leads to a set of linear equations for the unknown normal coordinate shifts  $\delta Q_s$ , and new coupling constants:

$$g_l = \pm \left( \frac{1}{\sqrt{2}} \right) \left[ \partial (\epsilon_k + F_k) / \partial Q_l \right]_0 \quad (12)$$

$$\gamma_{ll'} = \pm \left( \frac{1}{4} \right) \left[ \partial^2 (\epsilon_k + F_k) / \partial Q_l \partial Q_{l'} \right]_0$$

The coupling constants  $g_l, \gamma_{ll'}$  are calculated by the well-known perturbation expansion of the self-energy part. One could write:

$$\sum_{kk}^{(2)}(\epsilon) = \sum_{\substack{i,j \\ s \notin F}} \frac{(V_{ksij} - V_{ksji}) V_{ksij}}{\epsilon + \epsilon_s - \epsilon_i - \epsilon_j} + \sum_{\substack{i,j \\ s \notin F}} \frac{(V_{ksij} - V_{ksji}) V_{ksij}}{\epsilon + \epsilon_s - \epsilon_i - \epsilon_j} \quad (13)$$

and the coupling constant  $g_l$ , are as [17]:

$$g_l \approx \pm \frac{1}{\sqrt{2}} \frac{\partial \epsilon_k}{\partial Q_l} \frac{1 + q_k (\partial / \partial \epsilon) \sum_{kk} [-(V.I.P.)_k]}{1 - (\partial / \partial \epsilon) \sum_{kk} [-(V.I.P.)_k]} \quad (14)$$

The pole strength of the corresponding GF:

$$\rho_k = \left\{ 1 - \frac{\partial}{\partial \epsilon} \sum_{kk} [-(V.I.P.)_k] \right\}^{-1}; 1 \geq \rho_k \geq 0,$$

$$g_l \approx g_l^0 [\rho_k + q_k (\rho_k - 1)],$$

$$g_l^0 = \pm 2^{-1/2} \partial \epsilon_k / \partial Q_l \quad (15)$$

### 3. Fermi-liquid quasiparticle density functional theory

The quasiparticle Fermi-liquid version of the DFT [3-8,31,36] is used to determine the coupling constants etc. The master equations can be obtained on the basis of variational principle, if we start from a Lagrangian of a molecule  $L_q$ . It should be defined as a functional of quasiparticle densities:

$$v_0(r) = \sum_{\lambda} n_{\lambda} |\Phi_{\lambda}(r)|^2,$$

$$v_1(r) = \sum_{\lambda} n_{\lambda} |\nabla \Phi_{\lambda}(r)|^2, \quad (16)$$

$$v_2(r) = \sum_{\lambda} n_{\lambda} [\Phi_{\lambda}^* \Phi_{\lambda} - \Phi_{\lambda} \Phi_{\lambda}^*].$$

The densities  $v_0$  and  $v_1$  are similar to the HF electron density and kinetic energy density correspondingly; the density  $v_2$  has no an analog in the HF or DFT theory and appears as result of account for the energy dependence of the mass operator  $\Sigma$ . A Lagrangian  $L_q$  can be written as a sum of a free Lagrangian and Lagrangian of interaction:  $L_q = L_q^0 + L_q^{int}$ , where the interaction Lagrangian is defined in the form, which is characteristic for a standard DFT (as a sum of the Coulomb and exchange-correlation terms), but, it takes into account for a mass operator energy dependence of  $\Sigma$ :

$$L_q^{int} = L_K - \frac{1}{2} \sum_{i,k=0}^2 \int \beta_{ik} F(r_1, r_2) v_i(r_1) v_k(r_2) dr_1 dr_2 \quad (17)$$

where F is an effective exchange-correlation interaction potential. The constants  $\beta_{ik}$  are defined in Refs. [3-5]. The constant  $\beta_{02}$  can be calculated by analytical way, but it is very useful to keep in mind its connection with a spectroscopic factor  $F_{sp}$  [4,5]:

$$F_{sp} = \left\{ 1 - \frac{\partial}{\partial \epsilon} \sum_{kk} [-(V.I.P.)_k] \right\} \quad (18)$$

The new element is linked with using the DFT correlation Gunnarsson-Lundqvist, Lee-Yang-Parr functionals (c.g.[12-16]).

The multiplier  $[1 - \sum_2]$  is easily calculated if the Gunnarsson-Lundqvist-like correlation potential is used as  $V_{XC}$  and  $\sum_2$  is determined as follows:

$$\sum(r) = 0,254\rho(r)[0,328\rho^{-2/3}(r) + 0,204\rho^{-2/3}(r)/\{1+18,376\rho^{1/3}(r)\}]$$

#### 4. Results and conclusions

Let us present the results of calculating the binding energies and equilibrium distances in molecules  $N_2, O_2, F_2$  belonging to the class of complex from the point of view of taking into account the correlation effects. Effective approach to this topic can be performed with different versions of the standard DV approach such as: DV- $X_\alpha$ , DV- $X_\alpha$ (TS), MTO (see Table 1, where experimental data are also presented for comparison).

Table 1  
Bond energies (eV) and equilibrium distances (a.u.): 1 – experiment; 2- DV- $X_\alpha$ , 3- DV- $X_\alpha$ -TS, 4- MTO, 5 – Green function approach, 6 - present work

|   | $E_B$ | $R_e$ | $E_B$ | $R_e$ | $E_B$ | $R_e$ |
|---|-------|-------|-------|-------|-------|-------|
|   | $N_2$ |       | $O_2$ |       | $F_2$ |       |
| 1 | 9,91  | 2,07  | 5,22  | 2,28  | 1,67  | 2,68  |
| 2 | 9,39  | 2,11  | 8,15  | 2,35  | 8,97  | 2,35  |
| 3 | 10,86 | 2,11  | 4,14  | 2,33  | 3,54  | 2,61  |
| 4 | 7,8   | 2,16  | 4,10  | 2,30  | 0,6   | 2,91  |
| 5 | 9,99  | 2,11  | 3,78  | 2,33  | 3,21  | 2,6   |
| 6 | 9,95  | 2,09  | 5,12  | 2,31  | 1,58  | 2,70  |

As follows from the comparison of the energy values presented in Table 1 and the values of the molecular constants are sensitive to the calculation scheme. A more careful consideration of multiparticle correlation effects within the framework of this procedure leads to an improvement in the agreement between the calculation and experiment (in particular, this is observed for molecules; for there is some deterioration). The results obtained in this calculation are in slightly better agreement with experiment than the results of calculations in other versions of the DV method.

It can be concluded that the development of a more perfect than the existing methods

to computing molecular constants may be associated with careful consideration of complex correlation effects, including many-body corrections. The one-quasiparticle representation used here can be taken as the zeroth one in one of the sophisticated versions of the many-body PT such as , for example, the Möller-Plesset PT (see e.g. [6]).

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#### SPECTROSCOPIC FACTORS OF DIATOMIC MOLECULES: OPTIMIZED GREEN'S FUNCTIONS AND DENSITY FUNCTIONAL METHOD

**Summary.** It is presented an advanced approach to computing the energy and spectral parameters of the diatomic molecules, which is based on the hybrid combined density functional theory (DFT) and the Green's-functions (GF) approach. The Fermi-liquid quasiparticle version of the density functional theory is modified and used. The density of states, which describe the vibrational structure in photoelectron spectra, is defined with the use of combined DFT-GF approach and is well approximated by using only the first order coupling constants in the optimized one-quasiparticle approximation. Using the combined DFT-GF approach to computing the spectroscopic factors of diatomic molecules leads to significant simplification of the calculation procedure and increasing an accuracy of theoretical prediction. As illustration, the results of computing the bond energies in a number of known diatomic molecules are presented and compared with alternative theoretical results, obtained within discrete-variational  $X_\alpha$ , muffin-tin orbitals and other methods.

**Key words:** diatomic molecules, Green's functions, density functional

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#### ОПТИМИЗИРОВАННЫЙ МЕТОД ФУНКЦИОНАЛА ПЛОТНОСТИ И ФУНКЦИЙ ГРИНА В РАСЧЕТАХ ЭНЕРГИИ СВЯЗИ ДВУХАТОМНЫХ МОЛЕКУЛ

**Резюме.** Представлен усовершенствованный подход к вычислению энергетических и спектральных параметров двухатомных молекул, базирующийся на гибридной комбинированной теории функционала плотности (ТФП) и методе функций Грина (ФГ). Используется модель ферми-жидкостная квазичастичная версия ТФП. Плотность состояний, которая описывает колебательную структуру в фотоэлектронных спектрах, определяется с использованием комбинированного подхода ТФП - ФГ. Использование комбинированного ТФП-ФГ подхода приводит к значительному упрощению процедуры расчета и повышению точности теоретического прогнозирования. В качестве иллюстрации представлены результаты расчета энергий связи в ряде известных двухатомных молекул, которые сравниваются с альтернативными теоретическими результатами, полученными с помощью дискретно-вариационного  $X_\alpha$  и других подходов.

**Ключевые слова:** двухатомные молекулы, функция Грина, функционал плотности

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## **ОПТИМІЗОВАНИЙ МЕТОД ФУНКЦІОНАЛУ ГУСТИНИ І ФУНКЦІЙ ГРІНА В РОЗРАХУНКАХ ЕНЕРГІЇ ЗВ'ЯЗКУ ДВОАТОМНИХ МОЛЕКУЛ**

**Резюме.** Представлений вдосконалений метод обчислення енергетичних та спектральних параметрів 2-атомних молекул, що базується на гібридній теорії функціонала щільності (ТФП) і методі функцій Гріна (ФГ). Використано фермі-рідинну квазічастинкову версію ТФП. Густина станів, які описує коливальну структуру фотоелектронного спектру, визначається в межах ТФП-ФГ методу. Використання комбінованого ТФП-ФГ методу призводить до спрощення процедури обчислень, підвищення точності прогнозу. В якості ілюстрації представлені результати розрахунку енергій зв'язку в ряді відомих двоатомних молекул, які порівнюються з альтернативними теоретичними результатами, отриманими за допомогою дискретно-варіаційного і інших підходів.

**Ключові слова:** двоатомні молекули, функція Гріна, функціонал густини.